

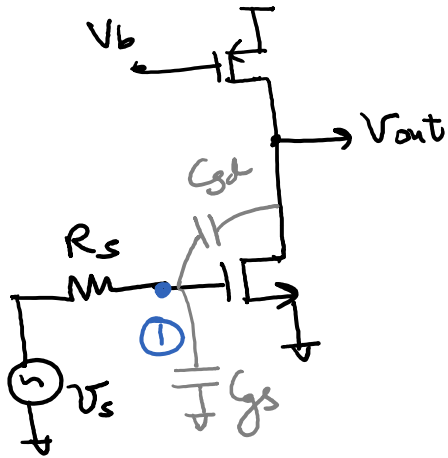
ECE SIS- Lecture 17

Thursday, October 18, 2018

11:01 AM

$$D(s) = \frac{s^2}{\omega_{p1}\omega_{p2}} + \left(\frac{1}{\omega_{p1}} + \frac{1}{\omega_{p2}}\right)s + 1$$

very small



$$\frac{v_{out}}{v_{in}}(s) = \frac{(sC_{gd} - g_{m1})R_o}{R_s R_o \zeta s^2 + [R_s(1 + g_{m1}R_o)C_{gd} + R_s C_{gs} + R_o(C_{gd} + C_o)]s + 1}$$

$$\zeta = C_{gs}C_{gd} + C_{gs}C_o + C_{gd}C_o$$

$$\omega_{p1} \Rightarrow \text{coeff of 's'}$$

$$= \frac{1}{R_s(1 + g_{m1}R_o)C_{gd} + R_s C_{gs} + R_o(C_{gd} + C_o)}$$

2nd pole: coeff of s^2 is $\frac{1}{\omega_{p1}\omega_{p2}}$

$$\omega_{p2} = \frac{1}{\omega_{p1}} \cdot \frac{1}{R_s R_o \zeta}$$

$$\Rightarrow \omega_{p2} = \frac{1}{\omega_{p1}} \times \frac{1}{R_s R_o [C_{gs1} C_{gd1} + C_{gs1} C_o + C_{gd1} C_o]}$$

$$= \frac{R_s (1 + g_m R_o) C_{gd1} + \boxed{R_s C_{gs1}} + R_o (C_{gd1} + C_o)}{R_s R_o [C_{gs1} C_{gd1} + C_{gs1} C_o + C_{gd1} C_o]}$$

* If $C_{gs1} \gg (1 + g_m R_o) C_{gd1} + \frac{R_o}{R_s} (C_{gd1} + C_o)$

$$\omega_{p2} \approx \frac{1}{R_o (C_{gd1} + C_o)}$$

← Same as ω_{out} in the Miller approximation

Valid only when C_{gs1} dominates the response, otherwise not.

Pole Splitting

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$$\omega_z = + \frac{g_{m1}}{C_{gd1}} \quad (\text{RHP zero})$$

$$\omega_{p1} = \frac{1}{R_s [(1 + |A_v|) C_c + C_{gs1}] + R_o (C_c + C_o)}$$

$$\omega_{p2} = \frac{R_s (1 + g_{m1} R_o) C_c + R_s C_{gs1} + R_o (C_c + C_o)}{R_s R_o [C_{gs1} C_c + C_{gs1} C_o + C_c C_o]}$$

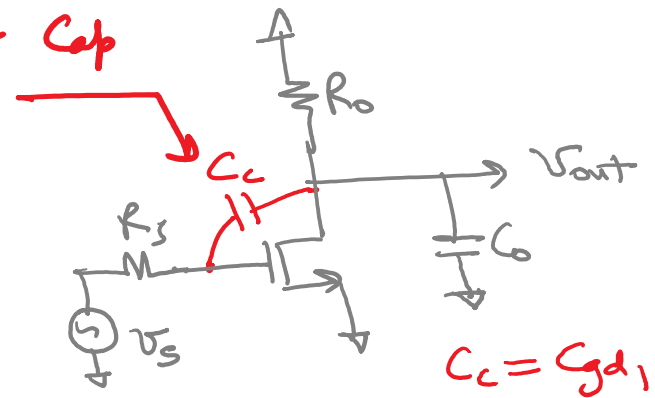
* for $C_c = 0$ and a large C_o

$$\omega_{p1} \approx \frac{1}{R_s C_{gs1} + R_o C_o}$$

$$\omega_{p2} \approx \frac{R_s C_{gs1} + R_o C_o}{R_s R_o C_{gs1} C_o} = \frac{1}{R_o C_o} + \frac{1}{R_s C_{gs1}}$$

$$\approx \frac{1}{R_o C_o}$$

Miller Cap



Now, increase C_c

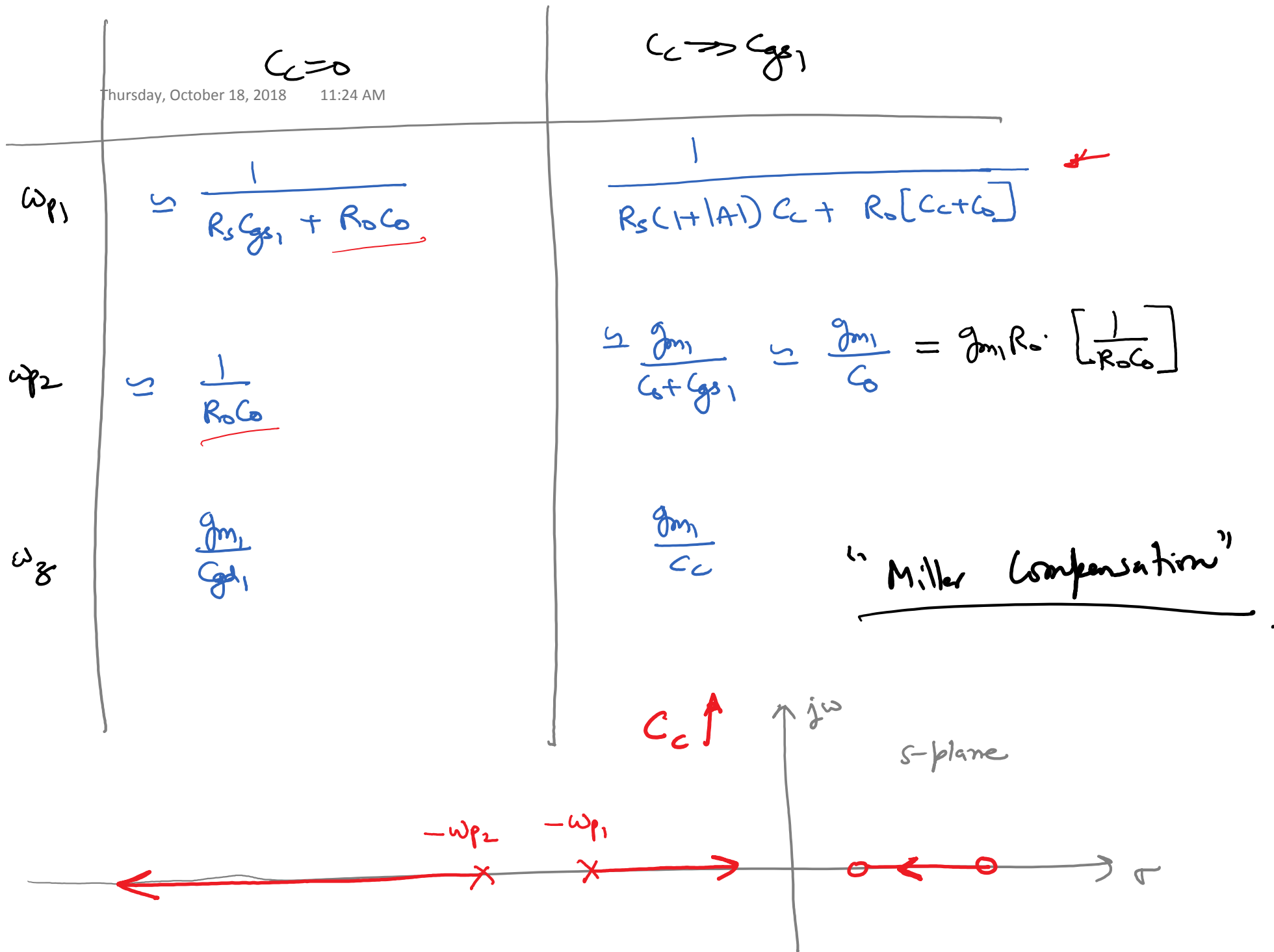
$$C_c \gg C_{gs1}$$

$$\omega_{p1} \approx \frac{1}{R_s (1 + |A_v|) C_c + R_o (C_c + C_o)}$$

$$\omega_{p2} \approx \frac{\cancel{R_s} (1 + \cancel{g_{m1} R_o}) \cancel{C_c} + \cancel{R_o} (C_c + C_o)}{\cancel{R_s R_o C_c} [C_o + C_{gs1}]}$$

very small

$$\approx \frac{g_{m1}}{C_o + C_{gs1}}$$



$$\omega_z = \frac{g_{m2}}{C_c}$$

RHP zero Cancellation



\Rightarrow



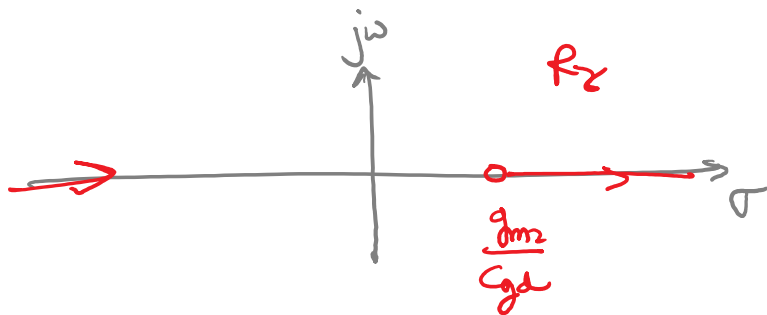
additional pole
at very high
frequencies

$$\omega_z = \frac{1}{C_c \left(\frac{1}{g_{m2}} - R_z \right)}$$

Set $R_z = \frac{1}{g_{m2}} \rightarrow$ zero is pushed to ∞
(it disappears)

$R_z > \frac{1}{g_{m2}} \rightarrow$ zero is pushed into the LHP
 \rightarrow LHP zero

\hookrightarrow add to the phase

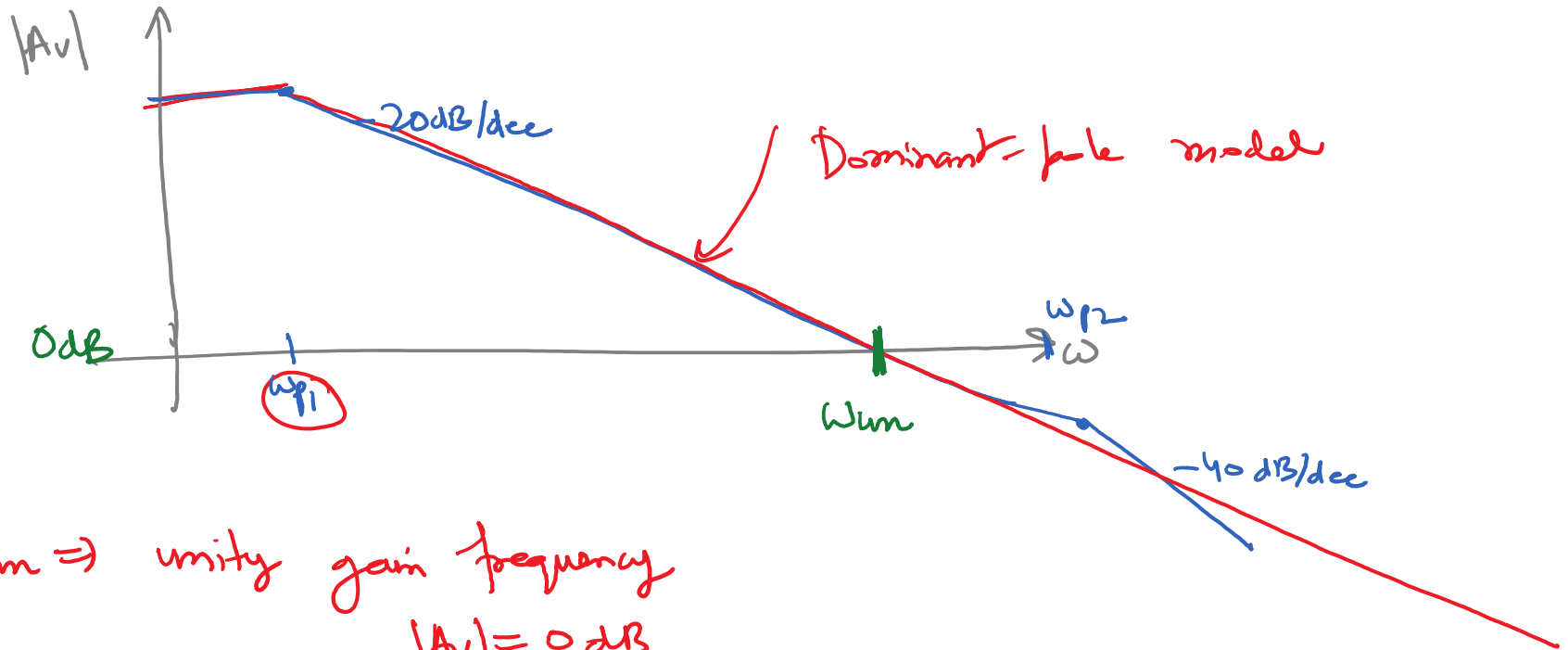


* Typically we just use R_z for
zero-nulling.

$$A_v(s) = \frac{A_{v0} (1 - s/\omega_z)}{(1 + s/\omega_{p1})(1 + s/\omega_{p2})} = \frac{A_{v0}}{(1 + s/\omega_{p1})(1 + s/\omega_{p2})}$$

$\omega_{p1} \ll \omega_{p2}$

zero nulling



$\omega_{um} \Rightarrow$ unity gain frequency
 $|A_v| = 0 \text{ dB}$

$\omega_{um} \Rightarrow$ depends upon $f_T \Rightarrow$

$f_{um} \approx \frac{f_T}{20}$ Empirical number

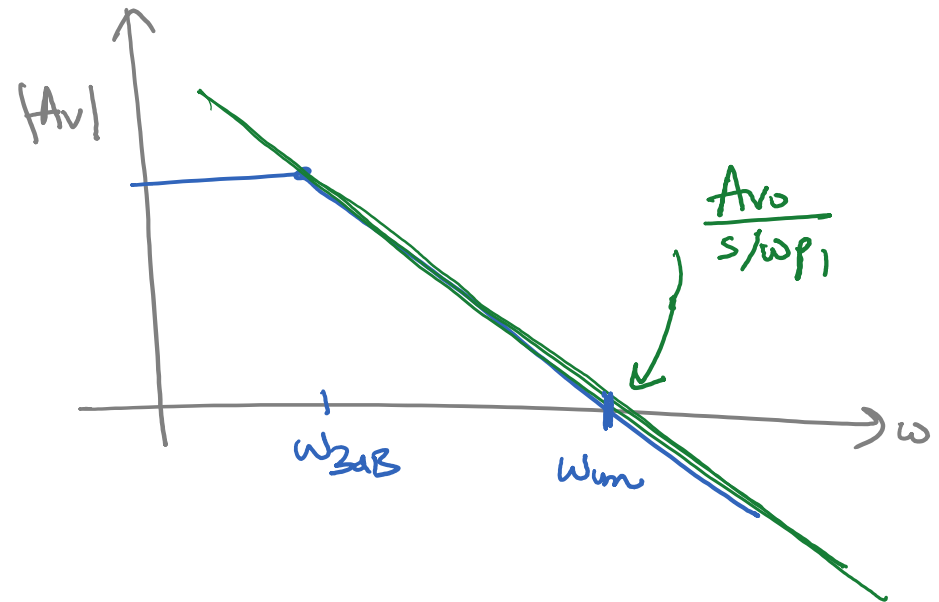
$$\omega_{3dB} \Rightarrow \omega_{p1}$$

Dominant pole Model

$$A_v(s) = \frac{A_{v0}}{(1 + s/\omega_{p1})} \approx \frac{A_{v0}}{(1 + s/\omega_{3dB})}$$

$$A_{v0} = g_{m1} R_1 g_{m2} R_2$$

$$\omega_{3dB} = \frac{1}{g_{m2} R_2 R_1 C_c}$$



Around $\omega \leq \omega_{um}$

$$A_v(s) \approx \frac{A_{v0}}{s/\omega_{3dB}}$$

$$A_v(s) \approx \frac{A_{v0} \omega_{3dB}}{s} \Rightarrow$$

$$A_v(j\omega_{um}) = \left| \frac{A_{v0} \omega_{3dB}}{j\omega_{um}} \right| = 1$$

$$\omega_{um} \approx A_{v0} \omega_{3dB}$$

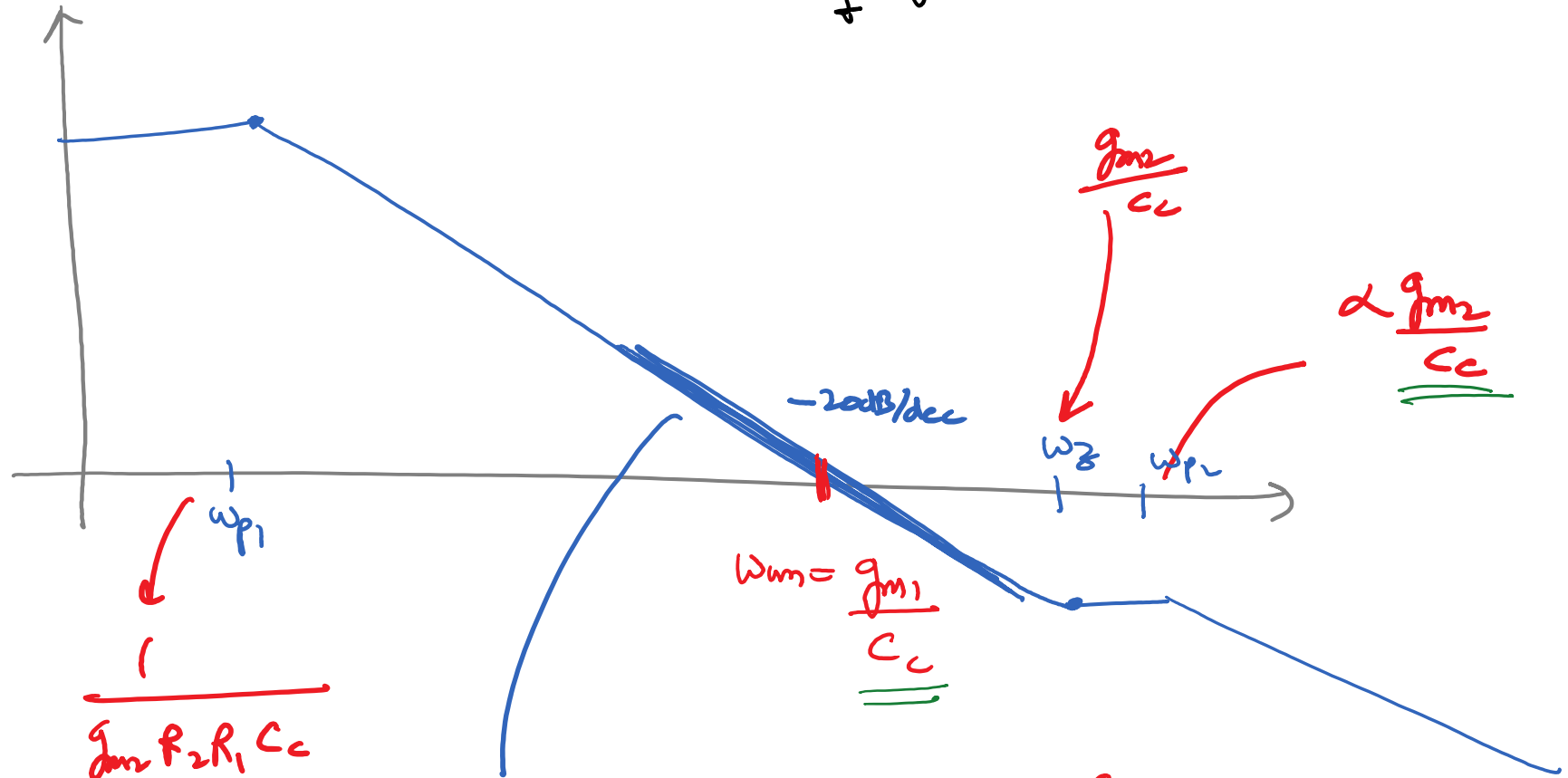
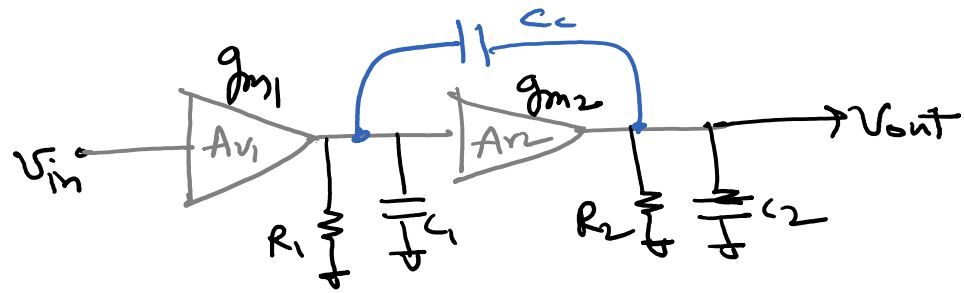
Gain-BW ←

$$W_{um} = A_{v0} \cdot W_{2dB}$$

$$= g_{m1} R_1 \cancel{g_{m2} R_2} \cdot \frac{1}{\cancel{g_{m2} R_2} R_1 C_c}$$

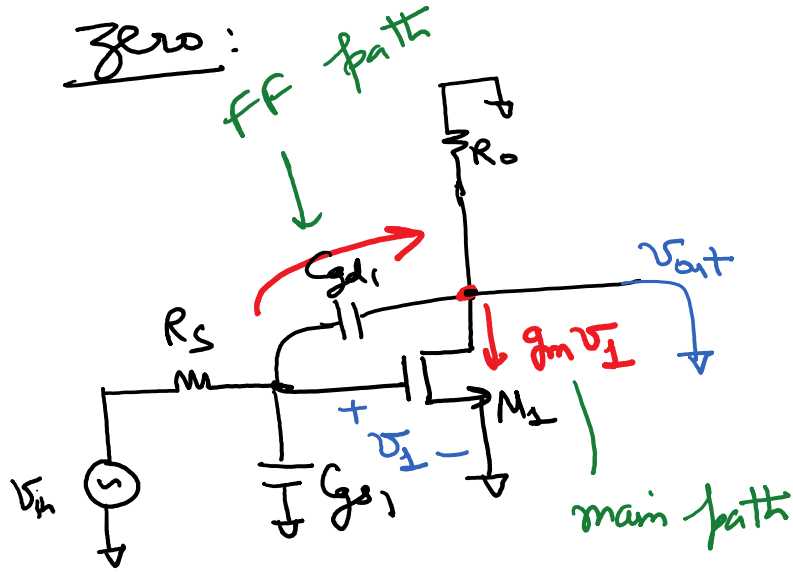
$$= \frac{g_{m1}}{C_c}$$

$$W_{um} = \frac{g_{m1}}{C_c}$$



* Should look like a first-order system around $\omega \leq \omega_{um}$ $g_{m2} > g_{m1}$

Zero:



* C_{gd1} provides a feedforward path that conducts input signal to the output at high frequencies.

→ adds to $g_m V_1$ in opposite phase

→ bring the phase down

→ RHP zero

At frequency

$\omega = \omega_z \Rightarrow$ these two currents cancel each other
→ gives rise to the zero.

* To estimate zero, output can be shorted to ground at $S = j\omega_z$

$$\therefore V_{out}(S = j\omega_z) = 0$$

\Rightarrow

$$V_1 \cdot C_{gd1} \cdot S_z = g_m V_1$$

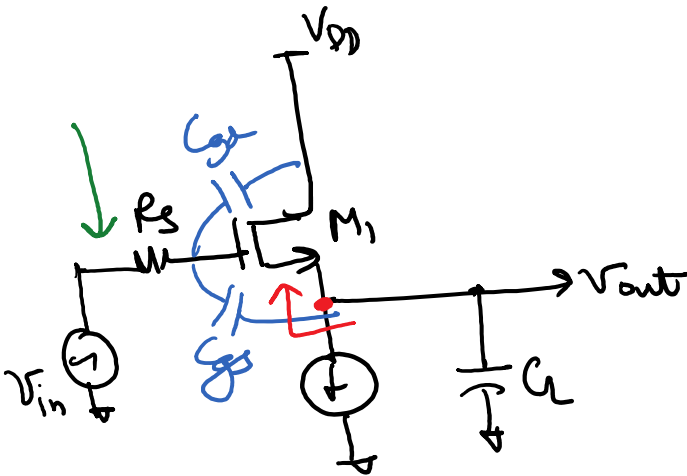
$$\Rightarrow S_z = \frac{g_m}{C_{gd1}}$$

$$\Rightarrow \omega_z = \frac{g_m}{C_{gd1}}$$

Source Follower Frequency Response:

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$$\frac{1}{R_C} = \frac{1}{\frac{1}{g_m} \cdot C_L} = \frac{g_m}{C_L}$$