

ECE 513- Lecture 16

Tuesday, October 9, 2018 3:28 PM

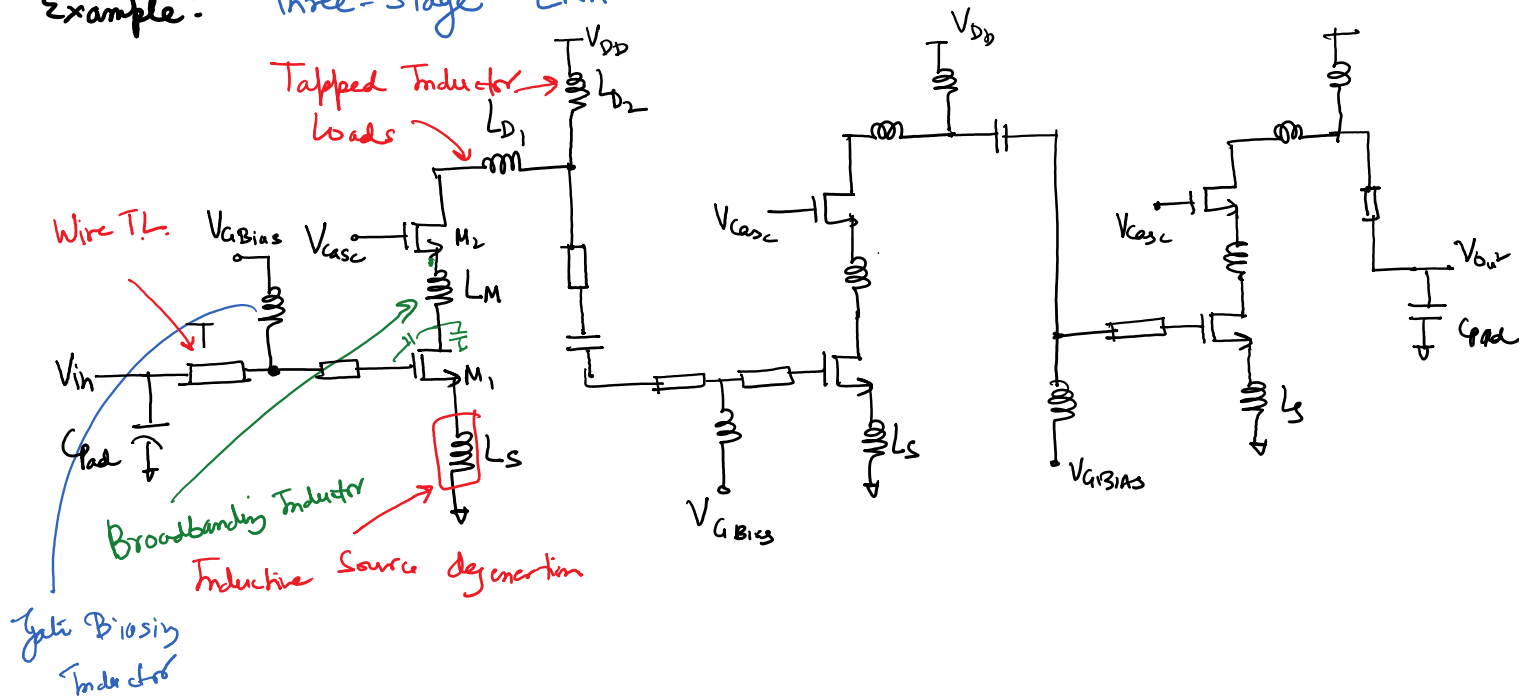
Chapter 5:

(T.H. Lee's Book)

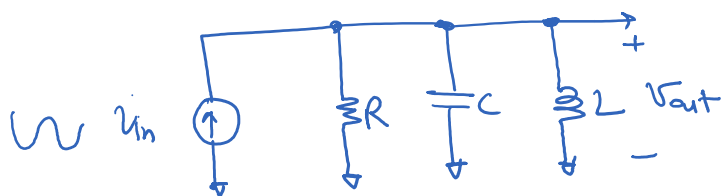
* A higher ratio of passive to active devices

(inductors, transformers, transmission lines)

Example: Three-stage LNA



Passive RLC Networks

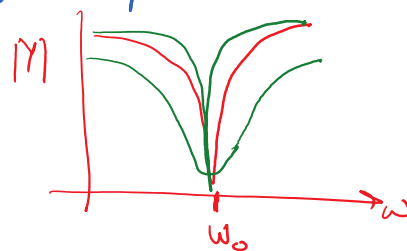


$$Y = g + j\omega C + \frac{1}{j\omega L} = \underset{\downarrow}{g} + j(\underset{\downarrow}{\omega C} - \frac{1}{\omega L})$$

Admittance goes to ∞ at DC \Rightarrow inductor is a short
 very high frequencies \Rightarrow capacitor is a short

At resonance frequency

$$\omega_0 C - \frac{1}{\omega_0 L} = 0 \Rightarrow \omega_0 = \frac{1}{\sqrt{LC}}$$



At resonance, admittance is purely real and equal to g

\hookrightarrow reactive terms cancel out

\hookrightarrow introduce Q

Q: quality factor

for a system is sinusoidal excitation

$$Q \triangleq \omega \cdot \frac{\text{Energy stored}}{\text{average power dissipated}} = 2\pi \cdot \frac{\text{Energy stored}}{\text{Energy dissipated in one cycle}}$$

Q → dimensionless

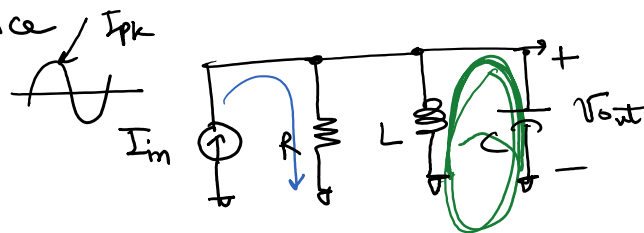
* If we neglect the loading of a network ⇒ unloaded-Q
including the load ⇒ loaded Q

↳ important to explicitly define the type of Q

Parallel RLC circuit at resonance

at resonance frequency, $\omega_0 = \frac{1}{\sqrt{LC}}$

$$V_{out} = |I_{in} R|$$



* Energy in LC sloshes back and forth with a constant sum

peak voltage across the capacitor = $I_{pk} R$

$$\Rightarrow \text{energy stored in the network} \Rightarrow E_{tot} = \frac{1}{2} C V_{pk}^2 = \frac{1}{2} C (I_{pk} R)^2$$

* Average power dissipated in the resistor

$$P_{avg} = \frac{1}{2} I_{pk}^2 R$$

* Q of the network at resonance

$$Q = \omega_0 \frac{E_{tot}}{P_{avg}} = \frac{1}{\sqrt{LC}} \cdot \frac{\frac{1}{2} C (I_{pk} R)^2}{\frac{1}{2} I_{pk}^2 R} = \frac{R}{\sqrt{L/C}}$$

* $\sqrt{L/C}$ has the dimensions of resistance

↳ characteristic impedance of the network

* Same as the magnitude of the capacitive and inductive reactances at resonance

$$|Z_C| = |Z_L| = \omega_0 L = \frac{L}{\sqrt{LC}} = \sqrt{\frac{L}{C}}$$

$$Q = \frac{R}{\sqrt{L/C}}$$

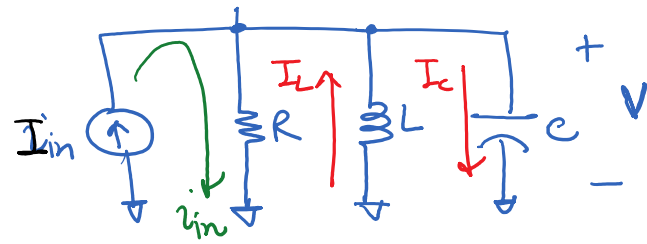
$$\text{as } R \rightarrow \infty \Rightarrow Q \rightarrow \infty$$

* Also,

$$Q = \frac{R}{|Z_{LC}|} = \frac{R}{\omega_0 L} = \omega_0 RC$$

$$Q = \frac{R}{\sqrt{4C}} = \frac{R}{\omega_0 L} = \omega_0 RC$$

Branch currents at resonance :



$$|I_C| = |I_L| = \frac{V}{Z} = \frac{|I_{in}|R}{\omega_0 L} = \frac{|I_{in}|R\sqrt{L/C}}{L}$$

$$= |I_{in}| \frac{R}{\sqrt{L/C}} = Q \cdot |I_{in}|$$

Ex. $Q = 1000$

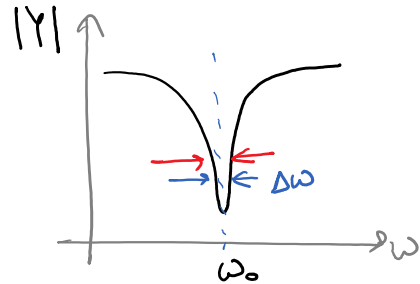
$I_{in} = 1 \text{ mA}$

$|I_L| = |I_C| = 1 \text{ A}$

Bandwidth & Q

Tuesday, October 9, 2018 4:11 PM

$$\omega = \omega_0 + \Delta\omega$$

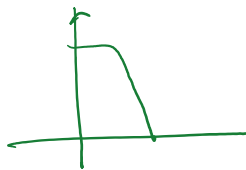
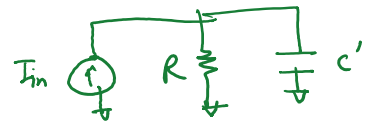


$$\begin{aligned} Y &= g + \frac{j}{\omega L} (\omega^2 LC - 1) = g + \frac{j}{\omega L} [(\omega_0 + \Delta\omega)^2 LC - 1] \\ &= g + \frac{j}{\omega L} \left[\frac{\omega_0^2 + \Delta\omega^2 + 2\omega_0\Delta\omega}{\omega_0^2} - 1 \right] \\ &= g + \frac{j}{\omega L} \left[\frac{\Delta\omega^2 + 2\omega_0\Delta\omega}{\omega_0^2} \right] \end{aligned}$$

for $\omega_0 \gg \Delta\omega$

$$\begin{aligned} Y &\approx g + \frac{j}{\omega_0 L} \cdot \frac{2\omega_0\Delta\omega}{\omega_0^2} \\ &= g + \frac{j2\Delta\omega}{\omega_0^2 \cdot L} \end{aligned}$$

$$Y = g + j2C\Delta\omega \equiv g + jC'$$



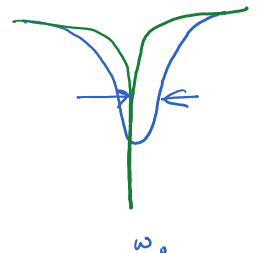
Define $BW = \frac{1}{RC}$

$$\Rightarrow \frac{BW}{\omega_0} = \frac{1}{RC\omega_0} = \frac{\sqrt{LC}}{RC} = \frac{\sqrt{L/C}}{R} = \frac{1}{Q}$$

fractional bandwidth

$$\frac{BW}{\omega_0} = \frac{1}{Q} \Rightarrow$$

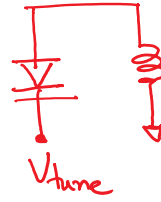
$$BW = \frac{\omega_0}{Q}$$



$$\omega_0 = \frac{1}{\sqrt{LC}}$$

$$= 1$$

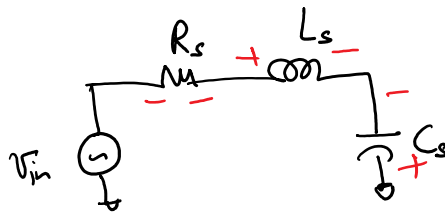




$$= \frac{1}{\sqrt{L(C_0 + \Delta C)}} \quad \text{VLC}$$

Series RLC Networks

- $\omega_0 = \frac{1}{\sqrt{L_s C_s}}$

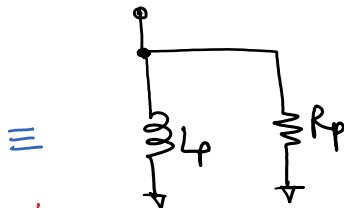
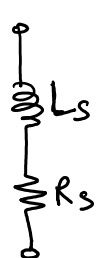


- $Q = \frac{\sqrt{L_s/C_s}}{R_s}$ ~~+~~ This is reciprocal of parallel RLC circuit.

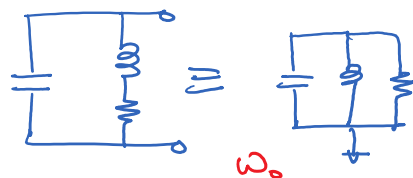
at resonance, ^{peak} voltage across either L or C is Q-times the voltage across the resistor

Series Parallel Transformation

Tuesday, October 9, 2018 4:27 PM



works only over a suitably small frequency range near resonance.



Let's equate impedance of series and parallel LR sections

$$j\omega_0 L_s + R_s = (j\omega_0 L_p) \parallel R_p = \frac{(\omega_0 L_p)^2 R_p + j\omega_0 L_p R_p^2}{R_p^2 + (\omega_0 L_p)^2}$$

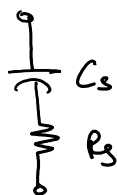
Equate real and imaginary parts

and use $Q = \frac{R_p}{\omega_0 L_p} = \frac{\omega_0 L_s}{R_s}$

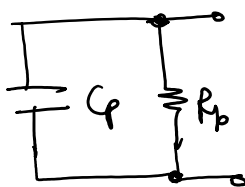
$$R_p = R_s (Q^2 + 1)$$

$$L_p = L_s \left(\frac{Q^2 + 1}{Q^2} \right)$$

Similarly.



\equiv



$$R_p = R_s (Q^2 + 1)$$

$$C_p = C_s \left(\frac{Q^2}{Q^2 + 1} \right)$$

Universal form

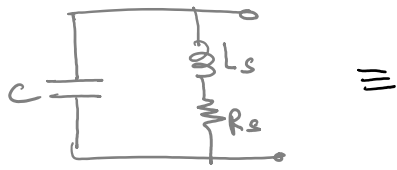
$$R_p = R_s (Q^2 + 1)$$

$$X_p = X_s \left(\frac{Q^2 + 1}{Q^2} \right)$$

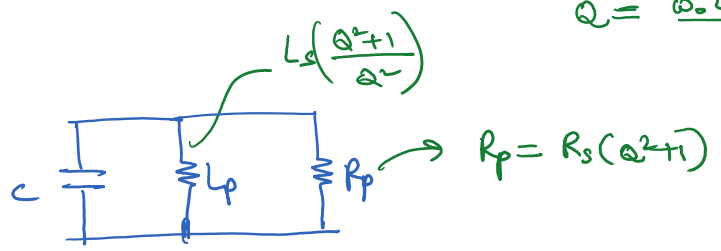
$$X = \frac{1}{\omega C} \text{ or } \omega L$$

* Can convert any "impure" RLC circuit into a purely parallel one which is straightforward to analyze.

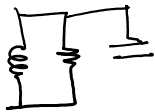
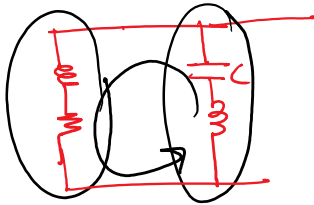
↳ But only valid for a narrow range of frequencies about ω_0 .



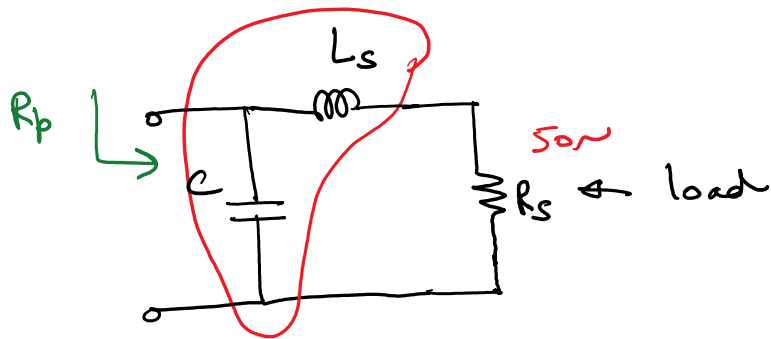
\equiv



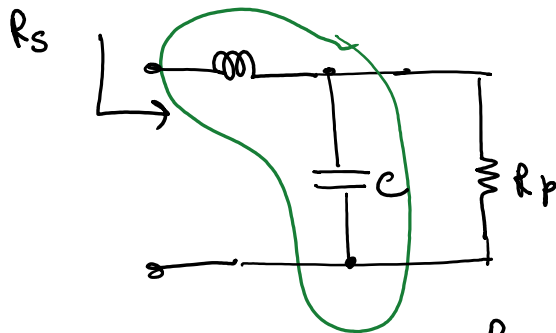
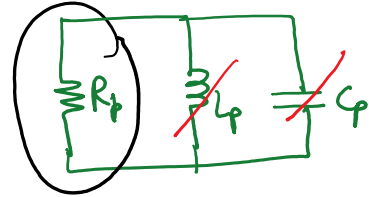
$$Q = \frac{\omega_0 L_s}{R_s} = \frac{R_p}{\omega_0 L_p}$$



L-match



$$R_p = (Q^2 + 1)R_s$$



$$R_s = \frac{R_p}{Q^2 + 1}$$

